

#### Introduction

Bra-Ket notation from Dirac

Measurements

Allowed transformations

Gates

1-bit full adder

Deutsch-Jozsa Problem

### Introduction to quantum computing dirac notation, measurements, unitary transformations, density matrix formalism

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2 Bra-Ket notation from Dirac









# Negative probabilities / 2-Norm[1]

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- Nature ist not described by probabilities but by *amplitudes* which can be positiv, **negative** or even complex
- Could have be invented by mathematicians
- Say  $(p_1 \dots p_N)$
- 1-Norm of  $(p_1 \dots p_N) = \sum_{i=0}^N p_i = 1$
- 2-Norm/Euclidean Norm is  $\sum_{i=0}^{N} p_i^2 = 1$



### Qubits

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- Deutsch-Jozsa Problem

- A quantumbit is 1 and 0 at the same time.
- By measuring it, the superposition collapses and the qubit is one of the measurementbases.
- In quantum theory is it common to write states with parenthesis like |·>.
- This notation is called Bra-Ket Notation or Dirac-Notation<sup>1</sup>

<sup>&</sup>lt;sup>1</sup>Physiciest Paul Dirac (1902 - 1984), Nobelprice together with Schroedinger 1933



### Qubits in Dirac-Notation

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### Qubit

A Quantumbit or Qubit is in state

$$\alpha \cdot |0\rangle + \beta \cdot |1\rangle$$

Where  $\alpha$  and  $\beta$  are *Amplitudes* and  $\alpha, \beta \in \mathbb{C}$  with

$$|\alpha|^2 + |\beta|^2 = 1$$



### Reminder: complex numbers

luction

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1-bit full adder	• Length $ a $ of a complex vector $a$ is $\sqrt{a}a$	

- Deutsch-Jozsa Problem
- where  $\overline{a}$  is the complex conjugate of  $a: a + ib \rightarrow a ib$



### Examples for Qubits

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• A qubit could be 
$$\frac{1}{\sqrt{2}} |0\rangle + \frac{1}{\sqrt{2}} |1\rangle$$

 $\blacksquare$  or like the *classical* bit  $0\cdot |0\rangle + 1\cdot |1\rangle = |1\rangle$ 



### Interpretation as vector space

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$$\binom{\alpha}{\beta} \in \mathbb{C}^2 = \alpha \binom{1}{0} + \beta \binom{0}{1} = \alpha \cdot |0\rangle + \beta \cdot |1\rangle$$

- Basis of the vector space is {|0⟩, |1⟩} so the superposition is getting a linearcombination of the basiselements.
- $\blacksquare$  Valid vectors must fulfill  $|\alpha|^2+|\beta|^2=1$



### Quantum registers

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Deutsch-Jozsa Problem  $\begin{array}{rcl} R & = & \left| x_1 \right\rangle \left| x_2 \right\rangle \\ & = & \left| x_1 x_2 \right\rangle \end{array}$ 

$$= (\beta_0 |0\rangle + \beta_1 |1\rangle) \cdot (\gamma_0 |0\rangle + \gamma_1 |1\rangle)$$

- $= \beta_0 \gamma_0 |00\rangle + \beta_0 \gamma_1 |01 + \beta_1 \gamma_0 |10\rangle + \beta_1 \gamma_1 |11\rangle\rangle$
- $= \alpha_{00} |00\rangle + \alpha_{01} |01\rangle + \alpha_{10} |10\rangle + \alpha_{11} |11\rangle$

■ aus 
$$|\beta_0|^2 + |\beta_1|^2 = 1$$
 und  $|\gamma_0|^2 + |\gamma_1|^2 = 1$ ) folgt[2]  
 $|\alpha_{00}|^2 + |\alpha_{01}|^2 + |\alpha_{10}|^2 + |\alpha_{11}|^2 = 1$ 



### Quantum registers

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- States of a quantum register with n bits are vectors in a 2<sup>n</sup> dimensional complex vector space
- Example for 2 bit:

$$|00\rangle = \begin{pmatrix} 1\\0\\0\\0 \end{pmatrix}, |01\rangle = \begin{pmatrix} 0\\1\\0\\0 \end{pmatrix}, |10\rangle = \begin{pmatrix} 0\\0\\1\\0 \end{pmatrix}, |11\rangle = \begin{pmatrix} 0\\0\\0\\1 \end{pmatrix},$$



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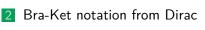
Measurement

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Measurements

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# Bra-Ket notation (Dirac)[3][4]

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- Kets like |0⟩ denote column vectors and are typically used to describe quantum states.
- $\{ \left| 0 \right\rangle, \left| 1 \right\rangle \}$  represent  $\{ (1,0)^T, (0,1)^T \}$
- **Bra**,  $\langle x |$  denotes the conjugate transpose of  $|x \rangle$ .
- Combining  $\langle x|$  and  $|y\rangle$  as in  $\langle x||y\rangle$ , also written as  $\langle x|y\rangle$ , denotes the **inner product** of two vectors.
- The notation  $|x\rangle \langle y|$  is the **outer product** of  $|x\rangle$  and  $\langle y|$ .



### Reminder: Dot product, scalar product

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$$\langle \vec{a}, \vec{b} \rangle = \vec{a} \cdot \vec{b} = \vec{b}^T \vec{a} = (a_0, a_1) \begin{pmatrix} b_0 \\ b_1 \end{pmatrix} = \sum_{i=0}^n a_i b_i$$

Inner product for Euclidian spaces



### Examples

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Deutsch-Jozsa Problem • Since  $|0\rangle$  is a unit vector  $\langle 0|0\rangle = 1$ 

• Since  $|0\rangle$  and  $|1\rangle$  are orthogonal we have  $\langle 0|1\rangle = 0$ 



# Outer product/Tensorproduct

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Deutsch-Jozsa Problem

$$\vec{u} \otimes \vec{v} = \vec{u}\vec{v}^{\mathsf{T}} = \begin{pmatrix} u_1 \\ u_2 \\ u_3 \\ u_4 \end{pmatrix} \begin{pmatrix} v_1 & v_2 & v_3 \end{pmatrix} = \begin{pmatrix} u_1v_1 & u_1v_2 & u_1v_3 \\ u_2v_1 & u_2v_2 & u_2v_3 \\ u_3v_1 & u_3v_2 & u_3v_3 \\ u_4v_1 & u_4v_2 & u_4v_3 \end{pmatrix}$$

■ Combining a *m*-dimensional vector with a *n*-dimensional vector results in a *m* × *n*-matrix

• 
$$|0\rangle \langle 1| = \begin{pmatrix} 1\\ 0 \end{pmatrix} (0,1) = \begin{pmatrix} 0 & 1\\ 0 & 0 \end{pmatrix}$$



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### Two types of operations

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- Deutsch-Jozsa Problem
- quantum state transformation



### Measuring

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### Measuring

If we measure a qubit in state  $\alpha \cdot |0\rangle + \beta \cdot |1\rangle$  the superposition is collapses (in another superposition). After measurement the qubit is with probability  $|\alpha|^2$  in state  $|0\rangle$  and with probability  $|\beta|^2$  in state  $|1\rangle$ .



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### Measurements

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### Conjugate transpose

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Deutsch-Jozsa Problem

$$A = (a_{ji}) \in \mathbb{C}^{m \times n}$$

$$A = \begin{pmatrix} a_{11} & \dots & a_{1n} \\ \vdots & \ddots & \vdots \\ a_{m1} & \dots & a_{mn} \end{pmatrix}$$

$$A^* = A^{\dagger} = \overline{A}^T = \overline{A}^T = \begin{pmatrix} \overline{a}_{11} & \dots & \overline{a}_{m1} \\ \vdots & \ddots & \vdots \\ \overline{a}_{1n} & \dots & \overline{a}_{mn} \end{pmatrix} \in \mathbb{C}^{n \times m}$$



# Unitary transformations

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Deutsch-Jozsa Problem •  $M^*$  or  $M^{\dagger}$  denotes the conjugate transpose / Hermitian transpose of the matrix M

Unitary operator

Matrix *M* is unitary if  $MM^* = M^*M = I$ 

- Unitary transformations are rotations or mirrorings in complex vector space
- Unitary transformations are reversible
- Unitary transformations are length-preserving  $||U|x\rangle || = |||x\rangle ||$
- For finite dimensional vector spaces  $M^*M = 1$  implies  $MM^* = 1$



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# Single qubit quantum state transformations / Pauli-gates

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Deutsch-Jozs Problem

$$\begin{array}{c|cccc} I: & |0\rangle \rightarrow |0\rangle & \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \\ \text{Pauli-X, Bitflip}: & |0\rangle \rightarrow |1\rangle & \begin{pmatrix} 0 & 1 \\ 1 \end{pmatrix} \\ \text{Pauli-Y}: & |0\rangle \rightarrow -|1\rangle & \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} \\ \text{Pauli-Y}: & |0\rangle \rightarrow -|1\rangle & \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix} \\ \text{Pauli-Z, Phaseflip}: & |0\rangle \rightarrow |0\rangle & \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \end{array}$$

1 - 1





# Hadamard<sup>2</sup> Transformation

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### Hadamard matrix

Matrix

Allowed trans-

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Deutsch-Jozsa Problem  $H = \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & 1\\ 1 & -1 \end{pmatrix}$ 

### is unitary and is called Hadamard Matrix

ŀ

$$\begin{aligned} \mathcal{H} : \quad |0\rangle &\to \frac{1}{\sqrt{2}} (|0\rangle + |1\rangle) \\ & |1\rangle &\to \frac{1}{\sqrt{2}} (|0\rangle - |1\rangle) \end{aligned}$$

<sup>2</sup>Mathematicians Jacques Hadamard (1865-1963)



# controlled-NOT gate $(C_{not})[5]$

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#### Deutsch-Jozsa Problem

 $\begin{array}{ccc} \mathcal{C}_{not}: & |00\rangle \rightarrow |00\rangle & \begin{pmatrix} 1 & 0 & 0 & 0 \\ |01\rangle \rightarrow |01\rangle & \\ |10\rangle \rightarrow |11\rangle & \\ |11\rangle \rightarrow |10\rangle & \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{pmatrix}$ 

- $C_{not}$  is unitary since  $C^*_{not} = C_{not}$  and  $C^*_{not}C_{not} = I$
- *C<sub>not</sub>* cannot be decomposed into a tensorproduct of two single-bit transformations





### Other gates

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Deutsch-Jozsa Problem

- Three bit controlled-controlled-NOT gate (Toffoli gate)<sup>34</sup>  $T = |0\rangle \langle 0| \otimes I \otimes I + |1\rangle \langle 1| \otimes C_{not}$
- Swap operation
  - $\mathbf{S}=\left|00\right\rangle \left\langle 00\right|+\left|01\right\rangle \left\langle 10\right|+\left|10\right\rangle \left\langle 01\right|+\left|11\right\rangle \left\langle 11\right|$
- Controlled swap (Fredkin gate)  $F = |0\rangle \langle 0| \otimes I \otimes I + |1\rangle \langle 1| \otimes S$

<sup>3</sup>Shi[6] had shown that the Hadamard and Toffoli gate already constitute a *universal set of quantum gates*.[1, p134]

 $^4$ Can be used to construct AND and NOT operators, 1-bit full adder



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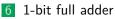
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# 1-bit full adder

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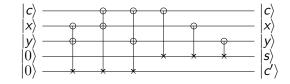
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with  $|x\rangle$  and  $|y\rangle$  being the databits,  $|s\rangle$  being the sum,  $|c'\rangle$  being the outcoming carrybit and  $|c\rangle$  being the incoming carrybit.

The Toffoli-gate is sufficient to construct arbitrary combinatorial circuit.



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### Deutsch-Jozsa Problem

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Gates	Deutsch-Jozsa Problem
1-bit full adder	Give a function $f: \{0, 1\} \rightarrow \{0, 1\}$ , say if f is constant
Deutsch-Jozsa Problem	$(f(0) = f(1))$ or balanced $(f(0) \neq f(1))$ . Running f is expensive.



# Classical solution

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Gates	Give a function $f: \{0, 1\} \rightarrow \{0, 1\}$ , say if f is constant
1-bit full adder	
Deutsch-Jozsa Problem	$(f(0) = f(1))$ or balanced $(f(0) \neq f(1))$ . Running f is expensive.

Running f two times and compare the results.



# Solution with quantum computing

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### Deutsch-Jozsa Problem

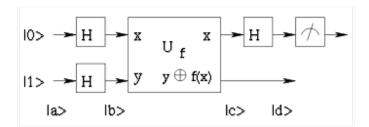
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Deutsch-Jozsa Problem Give a function  $f: \{0, 1\} \rightarrow \{0, 1\}$ , say if f is constant (f(0) = f(1)) or balanced  $(f(0) \neq f(1))$ . Running f is expensive.



Deutsch-algorithm, 1985



### Literatur I

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Problem

# Literatur II

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Gates	X Ch: "Deal as field and controlled and model little below."
1-bit full adder	F. Shi, "Both toffoli and controlled-not need little help to
Deutsch- Jozsa	do universal quantum computation,"